

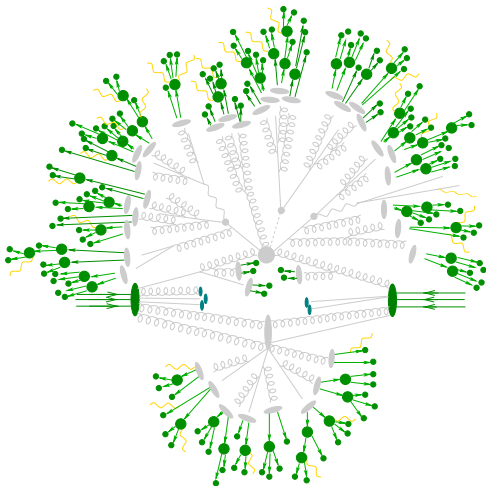


MC simulations of vector-boson production processes

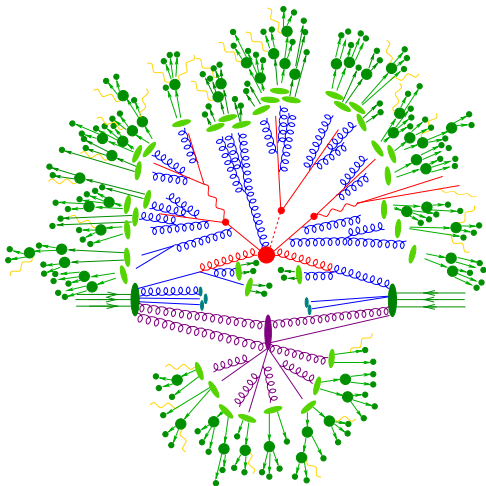
Recent methodological developments

Frank Siegert

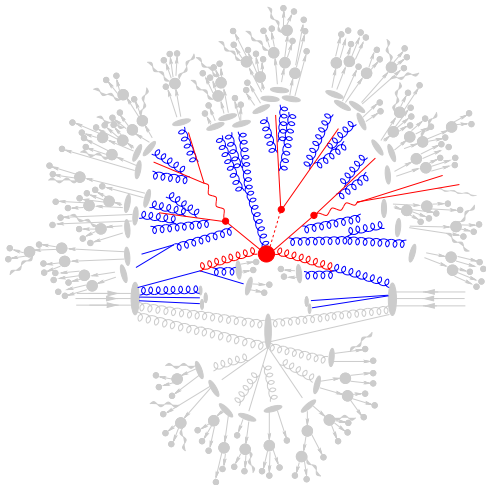
Standard Model at LHC 2014, Madrid, 9 April



- We want:
Simulation of $pp \rightarrow$ full hadronised final state
- MC event representation
- We know from first principles:
 - Hard scattering at fixed order in perturbation theory (**Matrix Element**)
 - Approximate resummation of QCD corrections to all orders (**Parton Shower**)
- Hadronisation/Underlying event



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Outline

- Improving parton showers at fixed order
 - (Applications in V+jets)
 - Applications in VV+jets
 - Finite loop² contributions
 - Applications in VVV+jets
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NLO+PS matching

- Parton shower on top of NLO prediction (e.g. inclusive W production)
- Objectives:
 - avoid double counting in real emission
 - preserve inclusive NLO accuracy



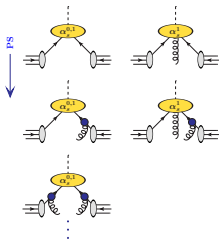
ME+PS@LO merging

- Multiple LO+PS simulations for processes of different jet multiplicity (e.g. W , W_j , W_{jj} , ...)
- Objectives:
 - combine into one inclusive sample by making them exclusive
 - preserve resummation accuracy



Combination: ME+PS@NLO

- Multiple NLO+PS simulations for processes of different jet multiplicity e.g. W , W_j , W_{jj} , ...
- Objectives:
 - combine into one inclusive sample
 - preserve NLO accuracy for jet observables



Basic idea

- “double-counting” between emission in real ME and parton shower
- ME is better than PS \rightarrow subtract PS contribution first
- but: shower unitary \rightarrow add “integrated” PS contribution back for NLO accuracy

Reminder + notation: NLO subtraction

$$d\sigma^{(\text{NLO})} = d\Phi_B \left[\mathcal{B} + \tilde{\mathcal{V}} + \sum_{\{ij\}} \mathcal{I}_{(ij)}^{(\text{S})} \right] + d\Phi_R \left[\mathcal{R} - \sum_{\{ij\}} \mathcal{D}_{ij}^{(\text{S})} \right]$$

NLO+PS formalism

- shower subtraction terms $\mathcal{D}_{ij}^{(\text{A})}$

$$d\sigma^{(\text{NLO sub})} = d\Phi_B \bar{\mathcal{B}}^{(\text{A})} + d\Phi_R \left[\mathcal{R} - \sum_{\{ij\}} \mathcal{D}_{ij}^{(\text{A})} \right]$$

$$\text{with } \bar{\mathcal{B}}^{(\text{A})} = \mathcal{B} + \tilde{\mathcal{V}} + \sum_{\{ij\}} \mathcal{I}_{(ij)}^{(\text{S})} + \sum_{\{ij\}} \int dt \left[\mathcal{D}_{ij}^{(\text{A})} - \mathcal{D}_{ij}^{(\text{S})} \right]$$

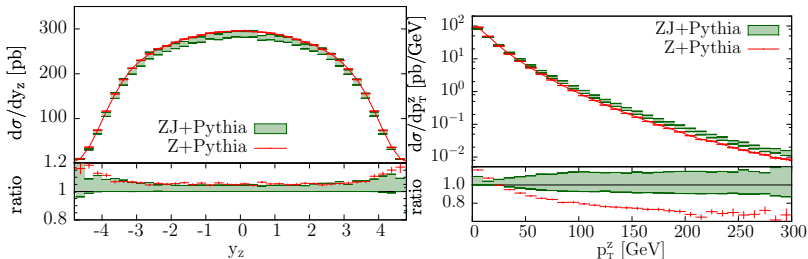
- apply PS resummation using $\mathcal{D}_{ij}^{(\text{A})}$ as splitting kernels

Master formula for NLO+PS up to first emission

$$\begin{aligned}
 d\sigma^{(\text{NLO+PS})} = & d\Phi_B \bar{\mathcal{B}}^{(A)} \left[\underbrace{\Delta^{(A)}(t_0, \mu_Q^2)}_{\text{unresolved}} + \underbrace{\sum_{\{ij\}} \int_{t_0}^{\mu_Q^2} dt \frac{\mathcal{D}_{ij}^{(A)}}{\mathcal{B}} \Delta^{(A)}(t, \mu_Q^2)}_{\text{resolved, singular}} \right] \\
 & + d\Phi_R \underbrace{\left[\mathcal{R} - \sum_{\{ij\}} \mathcal{D}_{ij}^{(A)} \right]}_{\text{resolved, non-singular} \equiv \mathcal{H}^{(A)}}
 \end{aligned}$$

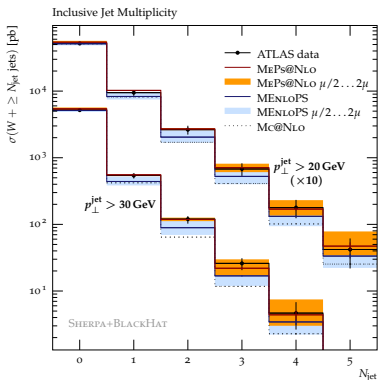
- To $\mathcal{O}(\alpha_s)$ this reproduces $d\sigma^{(\text{NLO})}$
- Event generation: $\bar{\mathcal{B}}^{(A)}$ or $\mathcal{H}^{(A)}$ seed event according to their XS
 - First line (“S-event”): from one-step PS with $\Delta^{(A)}$
 \Rightarrow emission (resolved, singular) or no emission (unresolved) above t_0
 - Second line (“H-event”): kept as-is \rightarrow resolved, non-singular term
- Resolved cases: Subsequent emissions can be generated by ordinary PS
- Exact choice of $\mathcal{D}_{ij}^{(A)}$ will specify MC@NLO vs. POWHEG vs. S-MC@NLO ...

Z+jet with MINLO

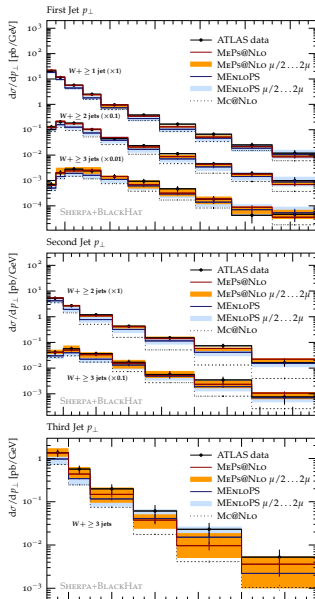


- POWHEG simulation for $pp \rightarrow Z+1$ jet, supplemented with CKKW-style scales and Sudakov form factors
 \Rightarrow inclusive simulation possible
- inclusive observables agree with $pp \rightarrow Z$ (within perturbative unc's)
- improved description of 1-jet observables: NLO for p_{\perp}^Z

W+jets with ME+Ps@NLO



- Comparison to ATLAS measurement
[Phys.Rev. D85 \(2012\), 092002](#)
- Significant reduction of ME+PS@NLO scale uncertainties in “NLO” multiplicities
- Improved agreement with data



Precise predictions for $pp \rightarrow ll\nu\nu + \text{jets}$

- **As signal:** SM measurements, vector-boson scattering, anomalous couplings, ...
- **As background:** Higgs production, BSM searches
- Higgs analyses in **exclusive 0, 1, 2-jet bins** (\Rightarrow jet vetoes)
 - \rightarrow Better control over backgrounds (WW^* vs. $t\bar{t}$)
 - \rightarrow Disentangle production modes ($gg \rightarrow H$ vs. VBF)

Non-trivial theoretical issues

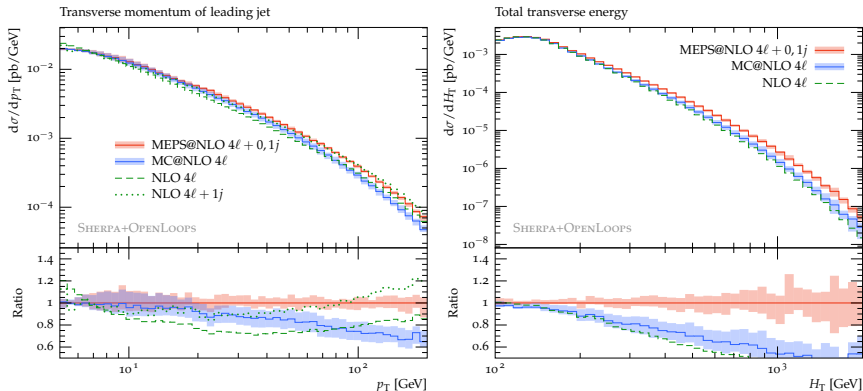
- Precise predictions for jet production \Rightarrow **beyond inclusive** NLO QCD
- Exclusive jet bins \Rightarrow Sudakov effects, **resummation**
- Offshell WW^* production \Rightarrow **non-resonant** and interference effects
- **Loop-induced** processes like $gg \rightarrow WW^*$ sizeable in Higgs signal regions

Toolkit

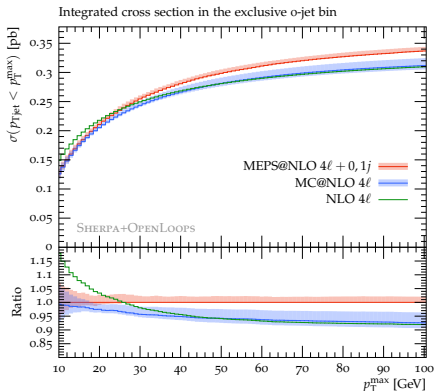
- SHERPA including its automated dipole subtraction and merging a la ME+PS@NLO
- OPENLOOPS automated 1-loop QCD matrix elements Cascioli, Maierhöfer, Pozzorini (2011)
- COLLIER for fast and stable tensor integral reduction Denner, Dittmaier, Hofer (in prep.)

\Rightarrow State-of-the-art QCD NLO automation

$p_{\perp, \ell} > 25 \text{ GeV}$, $|\eta_{\ell}| < 3.5$, $\cancel{E}_T > 25 \text{ GeV}$, anti- k_t jets with $R = 0.4$

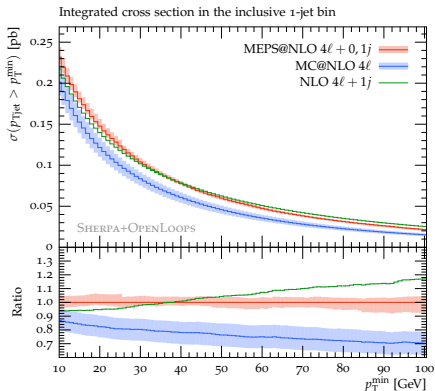


- NLO 4ℓ and S-MC@NLO 4ℓ only LO accurate, underestimate hard p_{\perp} tail
- Resummation necessary for $p_{\perp} \rightarrow 0$ (Sudakov logs)
 - NLO $4\ell \sim 20\%$ effects at $p_{\perp} = 5 \text{ GeV}$
 - NLO $4\ell + 1j$ partially includes logs \Rightarrow reduced effect
- Harder tails in fixed-order due to μ_R not dynamic with jet p_{\perp}
- H_T sensitive to combination of different jet multiplicities \Rightarrow merging crucial



Exclusive 0-jet bin

- Few-% agreement between S-Mc@NLO 4 ℓ and ME+PS@NLO
- Moderate Sudakov effects in comparison of NLO 4 ℓ and S-Mc@NLO 4 ℓ
- Low uncertainties \rightarrow good control wrt higher orders/logs

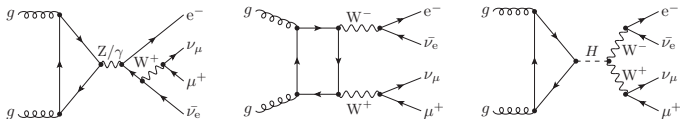


Inclusive 1-jet bin

- Sizable differences between S-Mc@NLO 4 ℓ and ME+PS@NLO, similar to jet p_{\perp}
- NLO 4 ℓ + 1j excess in tail due to α_s scale differences again

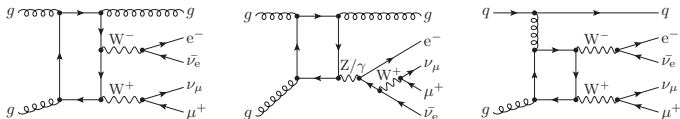
0-jet production: Examples for $gg \rightarrow 4\ell$ diagrams

- finite subset of NNLO contributions: squared quark loops like $gg \rightarrow 4\ell$
- relevant at LHC due to gluonic initial states, particularly in Higgs signal regions

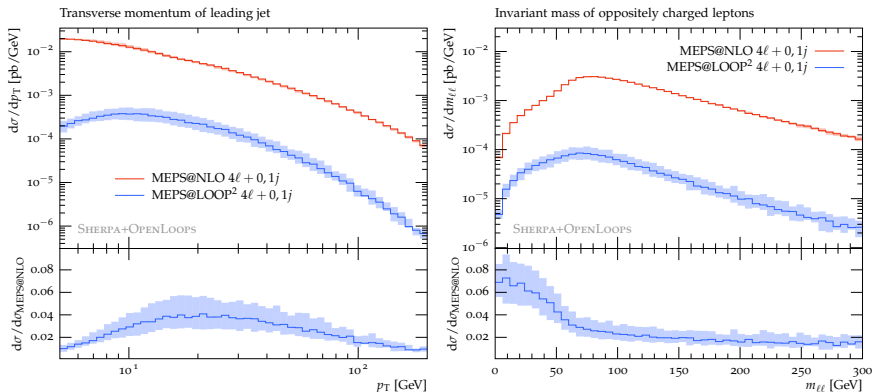


1-jet production

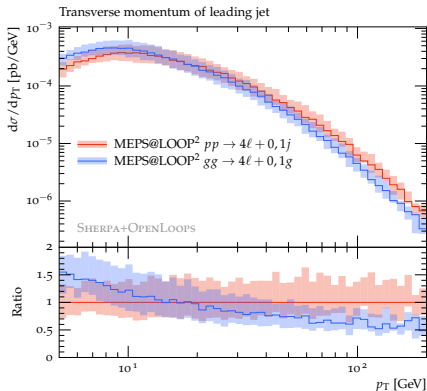
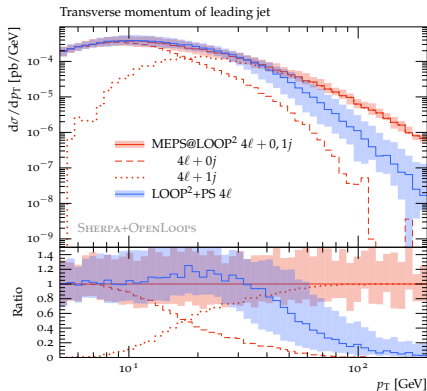
- example diagrams (requirement: vector bosons coupling to pure quark loop)



- first merging of 0-jet and 1-jet squared-loop contributions
- tree-level merging techniques since all MEs are finite
- shower on top of $gg \rightarrow 4\ell \Rightarrow$ consistency requires MEs for $qg, \bar{q}g$ and $q\bar{q}$ initial states



- Inclusive contribution of a few %
- Shape distortions: more significant impact in Higgs signal region (e.g. low $m_{\ell\ell}$)



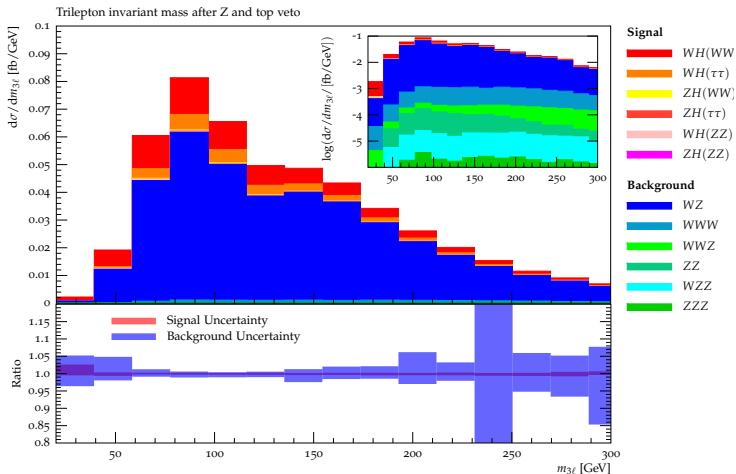
Merging effects

- Inclusion of LOOP² $4\ell + 1j$ in merging: harder p_{\perp} spectrum
- Significant reduction of uncertainties (wrt resummation scale) in high- p_{\perp} region

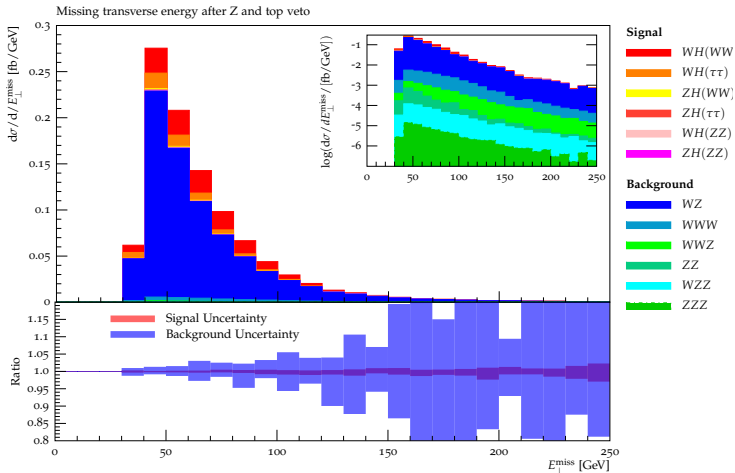
Non-gluonic initial states

- Inclusion of quark-channels → harder tail
- Naturally, lower Sudakov suppression without quark splittings
- Shape distortion
⇒ opposite effects in 0/1 jet bins

- trilepton analysis for associated Higgs production
- various Higgs signals and multi-boson backgrounds taken into account
- signals and dominant backgrounds (WZ, WWW) with ME+PS@NLO including +0,1 jets
 \Rightarrow uncertainty reduction to few-% level



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Summary

- recent conceptual developments in MC simulations of vector-boson production
- state of the art: multijet merging of NLO+PS matched simulations
 ~> cf. Stefan's talk tomorrow!
- application to single-, di- and triple-boson production:
 improved agreement with data, reduced uncertainties
- inclusion of loop²: finite, gauge-invariant part of NNLO contributions
 → shape distortions in 4ℓ production relevant for Higgs background
- Sherpa 2.x with ME+PS@NLO available publically

Thank you!

Backup material

$\text{Mc}\bar{\text{a}}\text{NLO}$

Frixione, Webber (2002)

$\mathcal{D}^{(A)} = \text{PS splitting kernels}$

- + Shower algorithm for Born-like events easy to implement
- "Non-singular" piece $\mathcal{R} - \sum_{ij} \mathcal{D}_{ij}^{(A)}$ is actually *singular*:
 - Collinear divergences subtracted by splitting kernels ✓
 - Remaining soft divergences as they appear in non-trivial processes at sub-leading N_c ✗

Workaround: \mathcal{G} -function dampens soft limit in non-singular piece
 \Leftrightarrow Loss of formal NLO accuracy
 (but heuristically only small impact)

$\text{S-Mc}\bar{\text{a}}\text{NLO}$

Höche, Krauss, Schönherr, FS (2011)

$\mathcal{D}^{(A)} = \text{Subtraction terms } \mathcal{D}^{(S)}$

- + "Non-singular" piece fully free of divergences
- Splitting kernels in shower algorithm become *negative*

Solution: Weighted $N_c = 3$ one-step PS based on subtraction terms

↓
Used in SHERPA

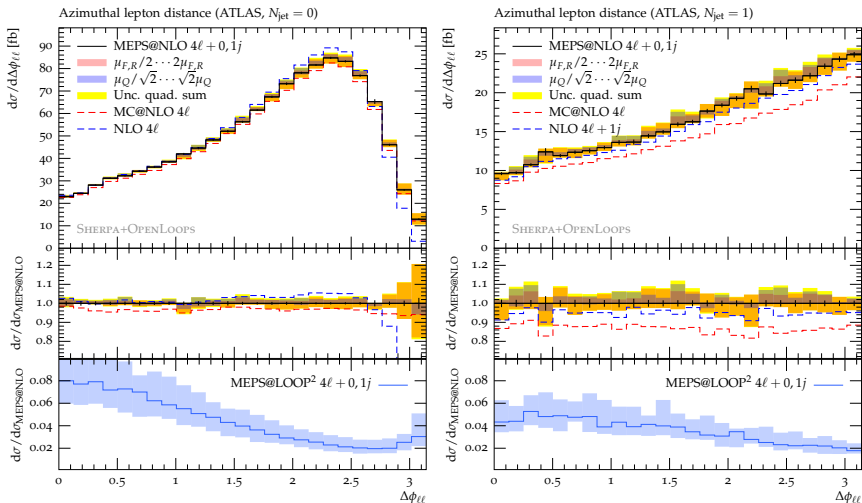
Rivet implementation of Higgs analyses

- 8 separate analyses: $\{\text{ATLAS,CMS}\} \times \{0\text{-jet, 1-jet}\} \times \{\text{signal region, control region}\}$
- Differential predictions in relevant observables: $p_{\perp}^j, m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T$

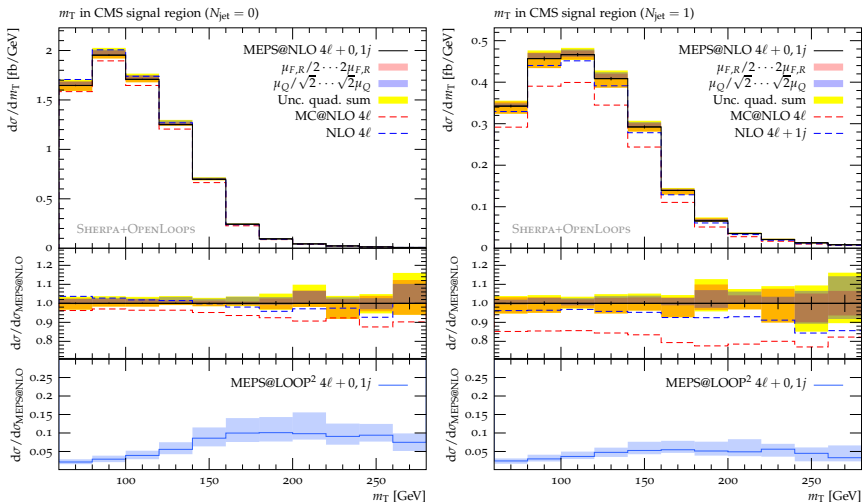
Findings

- Different simulation levels agree well in 0-jet bin (where they are NLO accurate)
- Fixed-order agrees with matched/merged predictions in most regions \rightarrow Sudakov logs not dominant, except e.g. $\Delta\phi_{\ell\ell} \rightarrow \pi$
- Pure MC@NLO predictions underestimates rate in 1-jet bins
- Uncertainty bands for best prediction (ME+PS@NLO) from $\mu_{R,F} \oplus \mu_Q$ variations at the few-% level

Example from ATLAS analysis



Example from CMS analysis



Signal/control cross sections in exclusive jet bins

- Relevant for background extrapolation from control to signal region in data-driven methods
- Example: ATLAS analysis

0-jet bin	NLO $4\ell (+1j)$	S-Mc@NLO 4ℓ	MEPS@NLO $4\ell + 0, 1j$	MEPS@LOOP ² $4\ell + 0, 1j$
σ_S [fb]	34.28(9) ^{+2.1%} _{-1.6%}	32.52(8) ^{+2.1%} _{-0.8%} ^{+1.2%} _{-0.7%}	33.81(12) ^{+1.4%} _{-2.2%} ^{+2.0%} _{-0.4%}	1.98(2) ^{+23%} _{-16.5%} ^{+27%} _{-20%}
σ_C [fb]	55.76(9) ^{+2.0%} _{-1.7%}	52.28(9) ^{+1.4%} _{-0.7%} ^{+1.4%} _{-1.1%}	54.18(15) ^{+1.4%} _{-1.9%} ^{+2.5%} _{-0.4%}	2.41(2) ^{+22%} _{-17%} ^{+27%} _{-18%}
1-jet bin	NLO $4\ell (+1j)$	S-Mc@NLO 4ℓ	MEPS@NLO $4\ell + 0, 1j$	MEPS@LOOP ² $4\ell + 0, 1j$
σ_S [fb]	8.99(4) ^{+4.9%} _{-9.5%}	8.02(4) ^{+8.5%} _{-6.4%} ^{+0%} _{-3.1%}	9.37(9) ^{+2.6%} _{-2.7%} ^{+2.5%} _{-0.1%}	0.46(1) ^{+40%} _{-18%} ^{+2.2%} _{-6.3%}
σ_C [fb]	26.50(8) ^{+6.4%} _{-12.5%}	24.58(8) ^{+6.1%} _{-6.5%} ^{+1.2%} _{-3.0%}	28.32(13) ^{+3.1%} _{-4.7%} ^{+4.1%} _{-0.0%}	0.79(1) ^{+33%} _{-20%} ^{+15%} _{-7%}

- Merged sample reproduces individual NLO cross sections well
- Combined uncertainty on ME+PS@NLO best prediction around 3(5)% in 0(1)-jet bin
- LOOP² effects larger in Signal than in Control region

Comparison of different simulation levels

NLO simulations	0-jet	1-jet	2-jet
NLO 4ℓ	NLO	LO	-
NLO $4\ell + 1j$	-	NLO	LO
S-Mc@NLO 4ℓ	NLO+PS	LO+PS	PS
S-Mc@NLO $4\ell + 1j$	-	NLO+PS	LO+PS
MEPS@NLO $4\ell + 0, 1j$	NLO+PS	NLO+PS	LO+PS
LOOP² simulations	0-jet	1-jet	2-jet
LOOP ² 4ℓ	LO	-	-
LOOP ² $4\ell + 1j$	-	LO	-
LOOP ² +PS 4ℓ	LO+PS	PS	PS
LOOP ² +PS $4\ell + 1j$	-	LO+PS	PS
MEPS@LOOP ² $4\ell + 0, 1j$	LO+PS	LO+PS	PS